

6 7 **Numerical Modeling of InGaAs/AlGaAs Heterostructure for Short-** 8 **Wave Infrared nBn Photodetectors**

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19 **Abstract** – nBn infrared photodetectors have emerged as a strong alternative to conventional pn and pin-based
20 structures due to their low dark current, fast response, and suppression of Shockley–Read–Hall (SRH) generation
21 pathways. In this study, an InP-based InGaAs/AlGaAs/InGaAs nBn heterostructure was designed and numerically
22 analyzed using SILVACO TCAD at 300 K to investigate its electro-optical behavior under various bias conditions.
23 The effects of barrier thickness and band alignment engineering were systematically evaluated with a particular focus
24 on minimizing the effective valence band offset and blocking majority-carrier leakage without hindering photocarrier
25 transport. The simulated band diagrams confirm that the AlGaAs barrier layer enables efficient majority-carrier
26 suppression while maintaining low-resistance conduction paths for photogenerated carriers. As a result, the proposed
27 device exhibits low dark current and high responsivity performance comparable to planar InGaAs structures even under
28 low bias operation. These results highlight the significance of optimized barrier design in achieving high detectivity
29 without relying on complex fabrication routes. Overall, the presented findings demonstrate the potential of tailored nBn
30 architectures for next-generation short-wave infrared imaging, low-photon-flux sensing, and high-dynamic-range
31 optoelectronic systems.

32 **Keywords:** nBn InGaAs, nBn SWIR, Low Noise, SILVACO TCAD, Low dark-current, InGaAs/AlGaAs
33 heterostructure

34 **1. Introduction**

35 Short-wavelength infrared (SWIR) photodetectors have become pivotal components in modern optoelectronic systems,
36 finding applications in defense imaging, medical diagnostics, environmental monitoring, and space-based instrumentation due
37 to their high quantum efficiency and broad spectral coverage [1-3]. Among these, InGaAs-based detectors stand out for their low
38 noise, lattice compatibility with InP, and high detectivity across the 0.9–1.7 μm range. However, conventional pn InGaAs
39 photodiodes increasingly fall short in meeting the demands of next-generation imaging systems that require ultra-low dark
40 current and stable operation at room temperature. In mesa-structured geometries, surface leakage and edge recombination effects
41 dominate, while in quadrant-type devices with large active areas, dark current rises sharply, limiting the achievable signal-to-
42 noise ratio [4-5].

43 To overcome these limitations, barrier-type architectures especially the nBn configuration have attracted significant research
44 attention [6-8]. The nBn concept, consisting of an n-type absorber, a wide-bandgap barrier, and an n-type contact layer, aims to

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block majority carrier (electron) flow while allowing minority carrier (hole) transport with minimal resistance. This configuration effectively suppresses surface and generation–recombination (G-R) dark currents without introducing a depletion region in the absorber [9]. A key design challenge, however, lies in the selection of an appropriate barrier material. Recent studies have explored a variety of barrier layers, including AlGaAs, InAlGaAs AlGaAsSb, and GaAsSb compounds, yet conflicting results have been reported regarding their valence-band alignment and hole transport efficiency [10-13]. These discrepancies underline the need for further investigation into band alignment optimization for nBn structures operating in the SWIR range.

In this study, we design and simulate an nBn InGaAs SWIR photodetector incorporating a lightly n-doped $\text{Al}_{0.57}\text{Ga}_{0.43}\text{As}$ barrier layer, which is hypothesized to provide near-flat valence-band alignment while maintaining sufficient conduction-band offset to block electron transport. Numerical simulations are performed using SILVACO ATLAS TCAD at 300 K, enabling accurate modeling of band profiles, carrier distributions, and current transport mechanisms. The optimized device exhibits a significant reduction in dark current compared to conventional pn InGaAs structures while preserving responsivity, indicating that proper band alignment through barrier engineering can lead to improved detectivity and thermal stability in SWIR photodetectors.

2. Simulation and Methodology

A detailed computational framework was developed to investigate the electrical transport and band-alignment characteristics of the proposed nBn-InGaAs photodetector. Simulations were performed using the SILVACO ATLAS, employing the physical models listed in Table 1 and the material parameters given in Table 2 [36-39].

Table 1. Physical Models used in SILVACO ATLAS simulation for nBn InGaAs photodetector design.

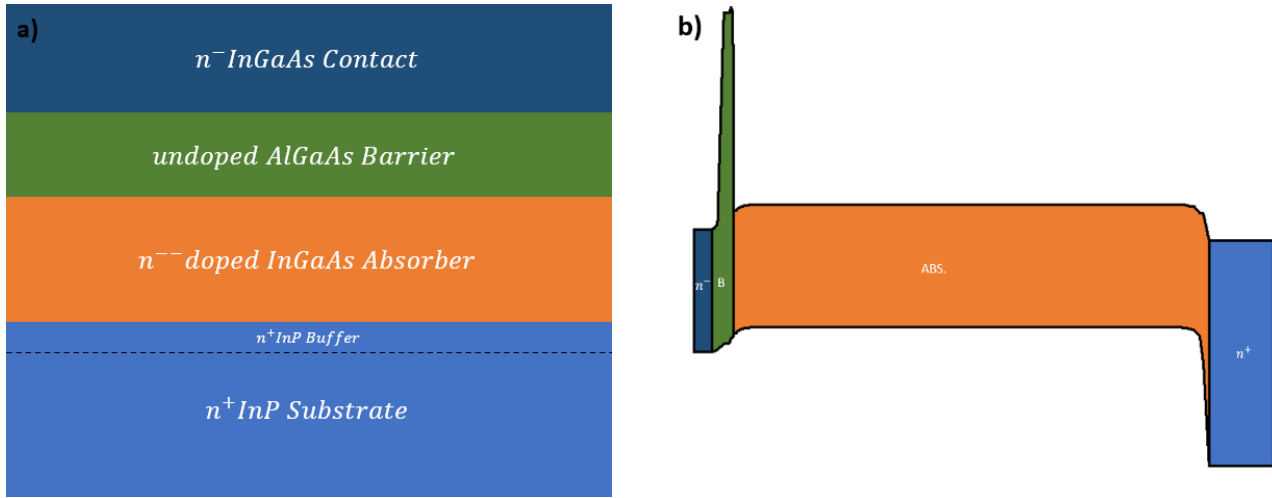
Model	Syntax	Description
Mobility	CVT	To enable transverse field, doping, and temperature-dependent parts of mobility
Optical	OPTR	To enable band-to-band recombination for direct band semiconductors
Shockley-Read-Hall	SRH	To enable recombination. It uses fixed minority carrier lifetimes
AUGER recombination	Auger	To enable direct transition of three carriers
Fermi distribution	Fermi	To enable carrier statistics. Suitable for highly doped regions
Quantum tunneling	QTUNN	To enable quantum tunneling through the conduction and valence band barrier due to a semiconductor
Quantum tunneling Band to Band	QTUN.BBT	To enables the band-to-band mode of the direct quantum tunneling model
Non-local trap assisted tunneling	TAT.NONLOCAL	To enables the non-local tunneling model in the calculation of the field effect enhancement factors.
Impact ionization	IMPACT SELB	To Enables the impact ionization. Recommended for most case, includes temperature dependent parameters.

In the simulation, the CVT mobility model was employed to account for doping, temperature, and transverse field effects on carrier mobility. Optical recombination was included via the OPTR model, while SRH and Auger models were used to describe trap-assisted and high-injection recombination processes, respectively. Fermi statistics were considered to ensure accurate carrier distribution in highly doped regions. Additionally, QTUNN and QTUN.BBT models were activated to capture quantum tunneling effects, including band-to-band tunneling across the barrier layer. To further improve the accuracy of dark current analysis, the TAT.NONLOCAL model was incorporated to account for non-local trap-assisted tunneling mechanisms, and the IMPACT SELB model was included to consider impact ionization effects with temperature dependence.

Table 2. Material parameters used in nBn InGaAs photodetector design.

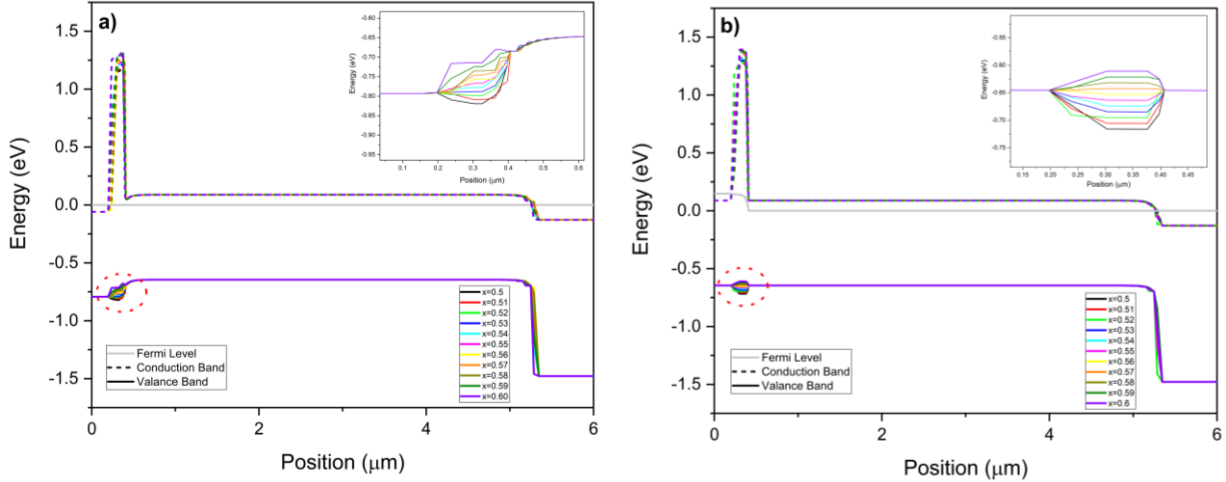
Parameter	InP	InGaAs	AlGaAs
Bandgap	1.35 eV	0.734 eV	1.99 eV
Electron affinity	4.40 eV	4.67 eV	3.41 eV
Permittivity	12.5	13.9	11.5
Electron SRH Lifetime	1 ns		
Hole SRH Lifetime	20 ns		
Trap Density	$1 \times 10^{14} \text{ cm}^{-3}$		
Electron Auger Coefficient	$3.7 \times 10^{-31} \text{ cm}^6/\text{s}$	$1.8 \times 10^{-28} \text{ cm}^6/\text{s}$	$5 \times 10^{-30} \text{ cm}^6/\text{s}$
Hole Auger Coefficient	$8.7 \times 10^{-30} \text{ cm}^6/\text{s}$	$1.8 \times 10^{-28} \text{ cm}^6/\text{s}$	$1 \times 10^{-31} \text{ cm}^6/\text{s}$
Electron Effective Mass	$0.0795 m_0$	$0.0332 m_0$	$0.77 m_0$
Hole Effective Mass	$0.6 m_0$	$0.471 m_0$	$0.734 m_0$

Vertical device stack (from substrate up) was modeled as InP (n^+ contact) / InGaAs (absorber) / $\text{Al}_x\text{Ga}_{1-x}\text{As}$ (barrier) / InGaAs (n^- contact) shown in Figure 1. Owing to its low lattice mismatch with the InGaAs and the advancements in epitaxial growth techniques, the AlGaAs barrier has emerged as the most suitable barrier material for InGaAs-based nBn photodetectors [14]. The reduced lattice strain at the heterointerface ensures high crystalline quality and stable electronic properties, facilitating efficient carrier transport while effectively suppressing dark current.

**Fig. 1.** (a) Planar design, and (b) band diagram of nBn structure.

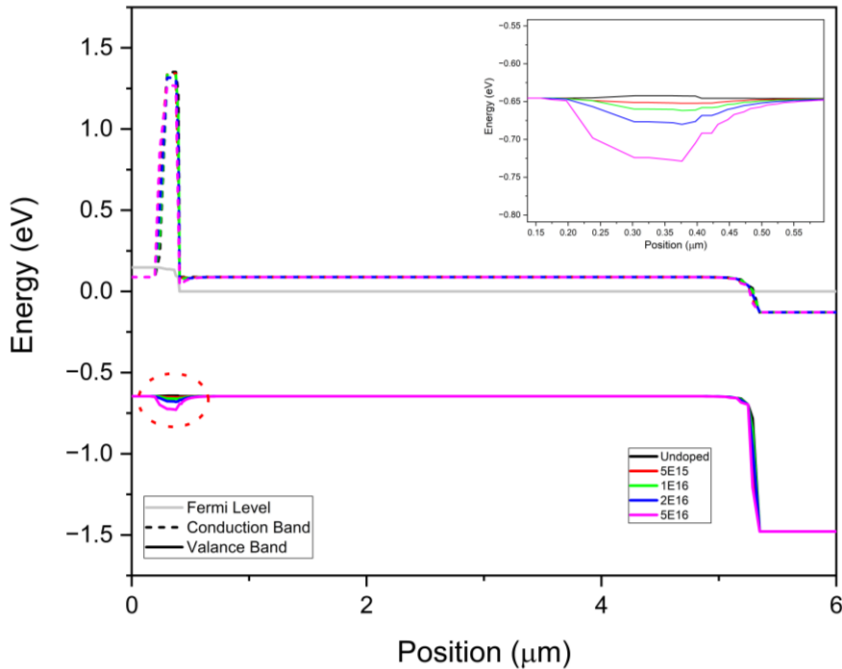
During the design of the nBn structure, systematic simulations were conducted to assess how different Al–Ga molar fractions ($\text{Al}_x\text{Ga}_{1-x}\text{As}$, with x ranging from 0.50 to 0.60) influence the electronic band alignment within the barrier layer. This compositional tuning was particularly crucial, as the barrier must balance two competing requirements: minimizing the valence band offset (VBO) to ensure unhindered hole transport, while simultaneously maintaining a sufficiently large conduction band offset (CBO) to suppress the flow of photogenerated electrons. Achieving this balance is essential for reducing dark current without compromising photocurrent extraction. The simulation results indicate that an Al composition of $x = 0.57$ provides the most favorable trade-off among the examined cases. At this composition, the VBO reaches its minimum value ($\Delta E_v \approx 35 \text{ meV}$), effectively eliminating the barrier for holes, whereas the CBO ($\Delta E_c \approx 1.26 \text{ eV}$) remains sufficiently high to block majority-carrier

87 leakage. Figure 2(a) illustrates the equilibrium band diagrams for the tested molar fractions, clearly demonstrating how
 88 incremental changes in Al content reshape both conduction and valence band edges. This optimized band configuration is
 89 particularly advantageous for nBn photodetectors, as it ensures that photogenerated minority carriers contribute efficiently to the
 90 photocurrent while majority carriers are strongly suppressed, ultimately leading to enhanced detectivity and improved bias
 91 stability.



92
 93 **Fig. 2.** Simulated energy band diagrams of the nBn InGaAs structure for different
 94 Al compositions : (a) under equilibrium conditions and (b) under -150 mV reverse bias voltage.

95 In Figure 2(b), the 150 mV reverse bias confines Fermi level band bending within the large bandgap AlGaAs barrier, leaving
 96 the InGaAs absorption layer largely undepleted and thus effectively mitigating G-R dark current, in line with the core design
 97 principles of nBn photodetectors [15]. Therefore, an Al composition of 0.57 was adopted for subsequent electrical and optical
 98 simulations to ensure optimal device performance.



99
 100 **Fig. 3.** Energy band diagrams of the nBn InGaAs structure under a 150 mV
 101 reverse bias for different N-type doping concentrations in the AlGaAs barrier.

102 Following the determination of the optimal Al composition ratio for the nBn InGaAs structure, the influence of varying N-type
 103 doping concentrations within the AlGaAs barrier layer was examined Figure 3 presents the corresponding energy band diagrams
 104 under a 150 mV reverse bias for different doping levels. As simulated, increasing the N-type doping concentration leads to a
 105 higher VBO, whereas the undoped AlGaAs barrier exhibits the most favorable band alignment. Nevertheless, to ensure that the
 106 simulation results accurately reflect realistic device behavior and considering that background doping cannot be completely
 107 avoided during epitaxial growth the electrical simulations were performed assuming a maximum donor concentration of 1×10^{16}
 108 cm^{-3} in the AlGaAs barrier. This consideration ensures a closer correspondence between the simulated and experimentally
 109 achievable device performance [16].

110 Based on the optimized band alignment parameters, the device structure was defined as follows: $0.2 \mu\text{m n}^-$ -In $_{0.53}$ Ga $_{0.47}$ As top n-
 111 contact layer ($5 \times 10^{17} \text{cm}^{-3}$), $0.1 \mu\text{m n}^-$ - Al $_{0.57}$ Ga $_{0.43}$ As barrier layer ($1 \times 10^{16} \text{cm}^{-3}$), $5 \mu\text{m n}^-$ - In $_{0.53}$ Ga $_{0.47}$ As absorber layer ($5 \times$
 112 10^{15}cm^{-3}), $0.5 \mu\text{m n}^+$ -InP bottom buffer n-contact layer ($5 \times 10^{18} \text{cm}^{-3}$) and $500\mu\text{m n}^+$ -InP substrate ($5 \times 10^{18} \text{cm}^{-3}$). The barrier
 113 thickness was determined according to the photocurrent versus potential curves presented in Figure 4 of Section 3. In addition,
 114 the absorber layer was designed to be relatively thick to achieve low capacitance, and lightly n-doped to further reduce the dark
 115 current [17].

116 3. Results and discussion

117 In this section, all simulations were conducted at 300 K, with the lateral device active area set to $300 \mu\text{m}$ to reflect the
 118 geometry of commercial detectors, and all current values were normalized accordingly. Figure 4(a) shows the simulated current
 119 as a function of applied bias for different barrier thicknesses. Increasing the barrier thickness clearly reduces the dark current,
 120 but at the same time causes a relative decrease in photocurrent. Detailed analysis revealed that a barrier thickness of 100 nm
 121 provides the optimal balance: it maintains photocurrent levels comparable to those of thinner barriers while yielding nearly an
 122 order-of-magnitude reduction in dark current. In contrast, barrier layers thicker than 100 nm introduce a more pronounced
 123 photocurrent degradation, indicating that additional thickness imposes unnecessary transport resistance without a commensurate
 124 benefit in dark-current suppression. Thus, to optimize the trade-off between dark-current reduction and photo-response, the
 125 barrier thickness in the nBn-InGaAs structure was selected as 100 nm [6]. In Figure 4(b), a sensitivity analysis was performed
 126 to evaluate the impact of band alignment uncertainties by varying the Al composition in the Al $_x$ Ga $_{1-x}$ As ($x = 0.55$ to 0.59) 100
 127 nm barrier. The results show that variations in Al content do not produce significant changes in dark current characteristics,
 128 indicating a weak dependence of device performance on moderate fluctuations in CBO and VBO. These findings confirm the
 129 robustness of the proposed nBn-InGaAs design against variations in material composition and the associated band offset
 130 uncertainties.

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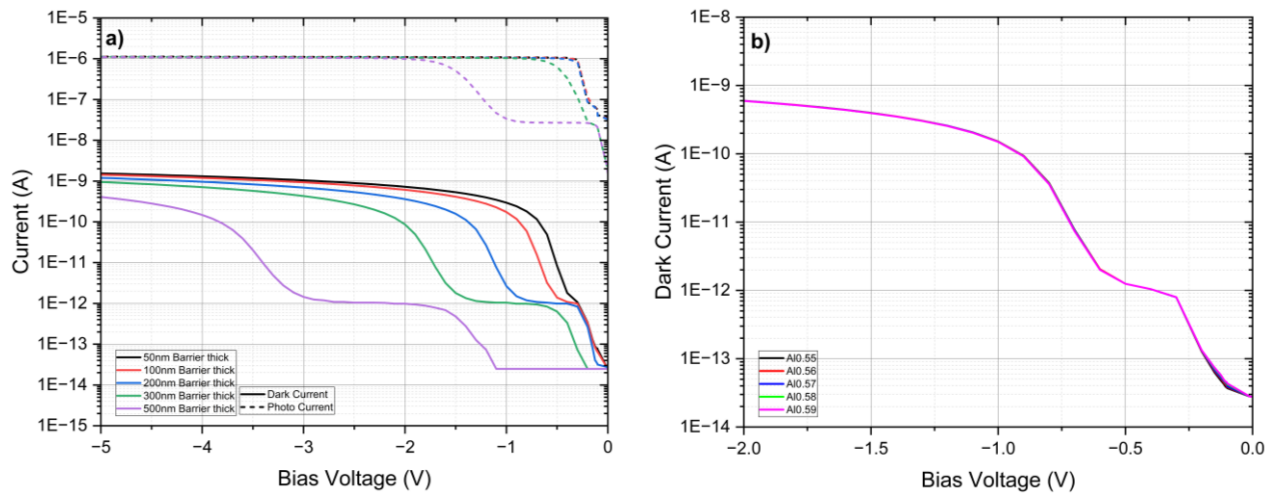
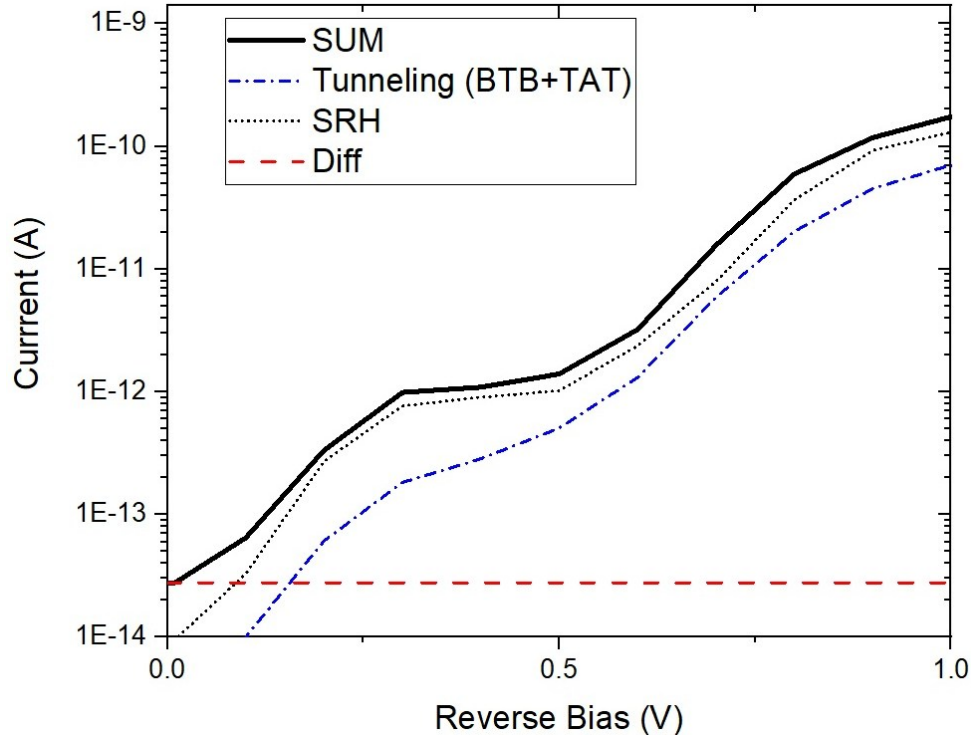


Fig. 4. (a) Dark and photocurrent of the nBn-InGaAs structure with various barrier thickness, (b) dark current graph with different Al ratio.

134 As illustrated in Figure 4(a), the nBn-InGaAs structure exhibits a significantly lower dark current compared to conventional
 135 mesa and planar InGaAs photodiodes [18-19], while maintaining nearly the same photocurrent magnitude. Furthermore, the
 136 photocurrent characteristics presented in Figure 4(a) were obtained under back-illuminated operation with an incident optical
 137 power density of 0.1 W/cm².
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Fig. 5. Dark current components as a function of reverse bias voltage at 300 K.

141 Figure 5 presents the decomposition of the dark current components as a function of reverse bias, providing insight into the
 142 dominant transport mechanisms in the proposed device. The total dark current (SUM) increases from approximately 3×10^{-14} A
 143 at 0 V to about 1×10^{-10} A at 1 V, exhibiting a gradual bias-dependent increase. Among the individual components, the SRH
 144 generation current dominates in the low-to-moderate bias range, indicating that generation–recombination processes are the
 145 primary contributor to the dark current [29-30].
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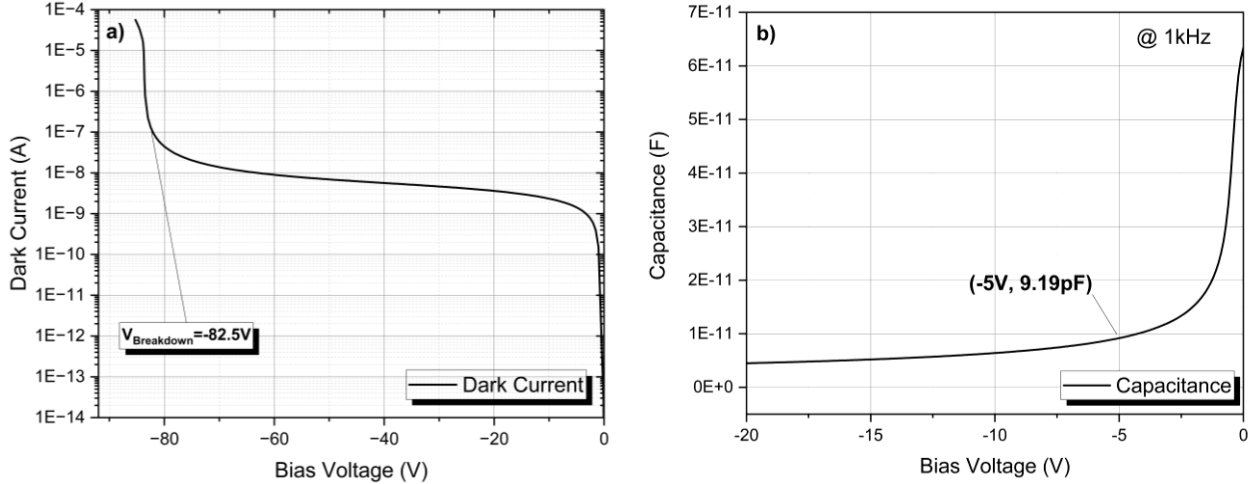
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147 In contrast, the tunneling current increases with reverse bias and becomes comparable to the SRH current at higher voltages,
 148 while remaining slightly lower overall, confirming that field-assisted leakage is effectively controlled by the barrier design. The
 149 diffusion current is nearly bias-independent and negligible in magnitude, suggesting minimal thermally driven carrier injection.
 150 Overall, these results demonstrate that the AlGaAs barrier effectively suppresses diffusion-related leakage, while the dark current
 151 is primarily governed by SRH generation at low bias and increasingly influenced by tunneling mechanisms at higher bias [31-
 152 32]. These results confirm that the optimized barrier design effectively minimizes leakage pathways without compromising
 153 carrier collection efficiency, demonstrating the superior dark-current suppression capability of the nBn structure.
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155 Following the analysis of the photocurrent characteristics, the simulated nBn-InGaAs structure was further evaluated in terms of
 156 its breakdown voltage and capacitance behavior, as presented in Figure 6. The results in Figure 6(a) show that the diode enters
 157 breakdown at approximately -82.5 V, primarily due to the combined effect of the thick absorber and barrier layers. This relatively
 158 high breakdown voltage highlights the structural robustness of the nBn configuration under strong reverse bias conditions and
 159 underscores its superior reliability compared to conventional pin InGaAs photodiodes [20], particularly in high-voltage and wide-
 160 dynamic-range applications. For comparison, conventional InGaAs pn photodiodes typically exhibit relatively low breakdown

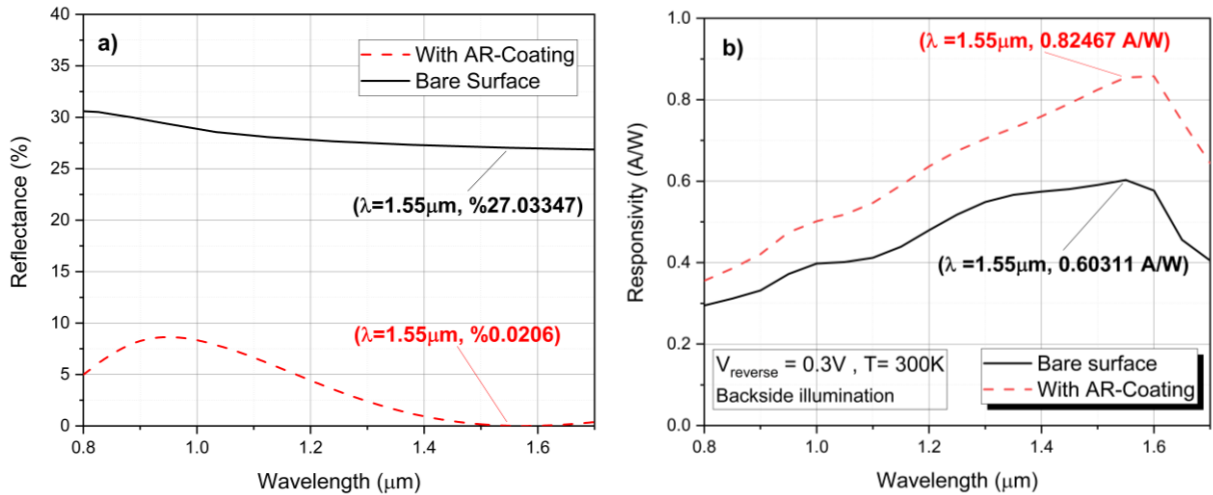
161 voltages due to the strong electric field concentration within the depletion region, which promotes early onset of avalanche
 162 multiplication through Impact Ionization. In such structures, the high doping gradients result in a narrow depletion width and
 163 consequently a large peak electric field, leading to breakdown at relatively low reverse biases. In contrast, nBn architectures
 164 eliminate the pn junction and instead rely on a unipolar barrier to control carrier transport. As a result, the electric field is more
 165 uniformly distributed across the device, significantly reducing the peak field intensity [27-28].



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Fig. 6. (a) Simulated I–V characteristics of the nBn-InGaAs structure and (b) C–V characteristics at 1 kHz.

168 **Figure 6(b)** shows that the device exhibits a capacitance of approximately 9.2 pF at –5 V bias and 1 kHz frequency, which can
 169 be attributed to the presence of the thick absorber layer. Such a low capacitance not only minimizes noise but also enhances
 170 temporal response and dynamic range, resulting in improved signal fidelity for high-speed and low-noise infrared detection [21].
 171 These findings demonstrate that, beyond its dark-current suppression capability, the nBn-InGaAs architecture combines
 172 electrical stability with excellent frequency response, confirming its potential for next-generation infrared photodetector
 173 technologies.



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Fig. 7. (a) Reflection and (b) responsivity graphs of the bare and AR-coated surfaces as a function of wavelength.

176 After evaluating the electrical characteristics, the photoresponsivity performance of the device was examined. Prior to
 177 responsivity analysis, the optical reflection behavior of the InP substrate was characterized, revealing that approximately 30%
 178 of the incident short-wavelength radiation is reflected from its surface. To enhance photon-to-electron conversion efficiency and
 179 improve overall responsivity, an anti-reflection (AR) coating was designed based on the reflectance formulation [22] in Eq (1).
 180 **Figure 7(a)** compares the reflectance spectra of a bare InP surface with structures incorporating Si₃N₄ (140 nm) and SiO₂ (150

181 nm) AR layers in the wavelength range of 0.8–1.7 μm . As illustrated, the AR-coated configuration significantly suppresses
 182 surface reflection, reducing it from nearly 30% to approximately 0.02% at 1.55 μm , which is expected to yield improved
 183 photocarrier generation and enhanced detector responsivity. Si_3N_4 and SiO_2 were selected as AR materials due to their low
 184 extinction coefficients in the sub-2 μm band, making them well-suited for near-infrared photodetector applications [23].

$$r_1 = \frac{n_0 - n_1}{n_0 + n_1} \quad (1)$$

$$r_2 = \frac{n_1 - n_2}{n_1 + n_2} \quad (2)$$

$$r_3 = \frac{n_2 - n_3}{n_2 + n_3} \quad (3)$$

$$\theta_1 = \frac{2\pi n_1 t_1}{\lambda} \quad (4)$$

$$\theta_2 = \frac{2\pi n_2 t_2}{\lambda} \quad (5)$$

$$R = \frac{r_1^2 + r_2^2 + r_3^2 + r_1^2 r_2^2 r_3^2 + 2r_1 r_2 (1 + r_3^2) \cos(2\theta_1) + 2r_2 r_3 (1 + r_1^2) \cos(2\theta_2) + 2r_1 r_2 r_3 \cos(2(\theta_1 + \theta_2)) + 2r_1 r_2 r_3 \cos(2(\theta_1 - \theta_2))}{1 + r_1^2 r_2^2 + r_2^2 r_3^2 + r_1^2 r_2^2 r_3^2 + 2r_1 r_2 (1 + r_3^2) \cos(2\theta_1) + 2r_2 r_3 (1 + r_1^2) \cos(2\theta_2) + 2r_1 r_2 r_3 \cos(2(\theta_1 + \theta_2)) + 2r_1 r_2 r_3 \cos(2(\theta_1 - \theta_2))} \quad (6)$$

185 To facilitate the definition of the reflectance expression in Eq.(6), the intermediate relations provided in Eq.(1-5) were utilized.
 186 Where n_0 is denotes the refractive index of the incident medium, n_1 the refractive index of the Si_3N_4 layer, n_2 the refractive index
 187 of the SiO_2 layer, and n_3 the refractive index of the InP substrate. Similarly, t_1 and t_2 represent the physical thicknesses of the
 188 Si_3N_4 and SiO_2 layers, respectively, while λ corresponds to the wavelength of the incident radiation.

189 In Figure 7(b), the wavelength-dependent responsivity characteristics of the nBn-InGaAs photodetector are presented. To clearly
 190 demonstrate the influence of the anti-reflection coating on photo-response, the responsivity curves of both the AR-coated device
 191 and the bare surface are shown together. As observed, the application of the AR layer yields a substantial enhancement in
 192 photocarrier extraction efficiency: at the telecommunication wavelength of 1.55 μm [24], the responsivity increases by
 193 approximately 36%, improving from 0.60 A/W to 0.82 A/W. This improvement is consistent with the reflectance reduction
 194 achieved through AR design and confirms that minimizing front-surface optical losses directly enhances the detector's photon-
 195 to-electron conversion capability. In the numerical simulations performed in this study, the TCAD software provides external
 196 quantum efficiency (EQE) rather than responsivity as a direct output. Therefore, the EQE values were converted to responsivity
 197 using Eq. (7), enabling a direct comparison between the simulated and AR-enhanced device performance [25]. It should also
 198 be noted that the reported responsivity values correspond to conditions in which the nBn-InGaAs structure is illuminated from
 199 the backside under an incident optical power density of 0.01 W/cm². Backside illumination ensures that the incident photons
 200 efficiently reach the absorbing InGaAs layer without encountering additional front-side metallization losses, thereby allowing a
 201 more accurate assessment of the intrinsic photo-response. Collectively, these results verify that integrating an appropriately
 202 engineered AR coating is an effective approach for maximizing optical coupling and elevating responsivity in nBn-type infrared
 203 photodetectors.

$$R = \eta \frac{q}{hf} \quad (7)$$

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205 where η is the external quantum efficiency of the detector for a given wavelength, q is the electron charge, f is the frequency of
 206 the incident photons, and h is the Planck constant.

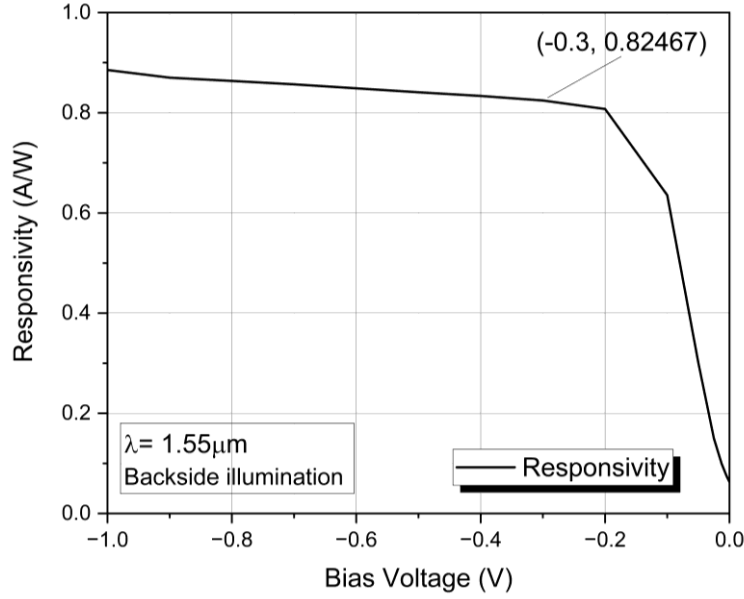


Fig. 8. Bias voltage-dependent responsivity graph of nBn-InGaAs structure.

Figure 8 shows the responsivity variation with applied bias voltage, indicating that the proposed nBn design maintains high responsivity even at very low operating bias due to the minimized VBO and optimized barrier engineering. Notably, its performance approaches that of cutting-edge planar InGaAs detectors (Table 3), confirming the viability of the structure as a strong candidate for high-performance, low-power infrared photodetection.

Table 3. Performance comparison of the proposed nBn photodetector with reported InGaAs-based photodetector

Device	Dark Current	Responsivity	EQE	Ref
nBn-InGaAs	~1 pA ($V_r = 0.3$ V, $\phi = 300\mu\text{m}$)	0.824 A/W ($\lambda = 1.55$ μm , $V_r = 0.3$ V)	-	This work
nBn-InGaAs	~8 nA ($V_r = 0.3$ V, $\phi = 300\mu\text{m}$)	-	% 45 ($\lambda = 1.55$ μm , $V_r = 0.25$ V)	[37]
nBn-InGaAs	~0.9 nA ($V_r = 0.3$ V $\phi = 300\mu\text{m}$)	-	% 80 ($\lambda = 1.55$ μm , $V_r = 0.2$ V)	[6]
pBin-InGaAs	~70 pA ($V_r = 0.3$ V $\phi = 500\mu\text{m}$)	0.62 A/W ($\lambda = 1.55$ μm , $V_r = 0.2$ V)	-	[33]
pn-InGaAs	~10 pA ($V_r = 0.3$ V $\phi = 300\mu\text{m}$)	1.08 A/W ($\lambda = 1.55$ μm)	% 86 ($\lambda = 1.55$ μm)	[34]
pin-InGaAs	~41 pA ($V_r = 5$ V $\phi = 73\mu\text{m}$)	0.99 A/W ($\lambda = 1.55$ μm $V_r = 5$ V)	% 79 ($\lambda = 1.55$ μm)	[35]

To benchmark the proposed device, a comparison with previously reported InGaAs photodetectors is summarized in Table 3. The proposed nBn structure exhibits an ultra-low dark current of 1 pA at -0.3 V, which is significantly lower than reported nBn-InGaAs devices (0.9–8 nA) under similar bias conditions, corresponding to an improvement of several orders of magnitude. Compared to conventional pn junction, PIN diode, and pBin structures, which typically show dark currents in the range of 10–70 pA, the proposed device still demonstrates superior suppression.

The responsivity of 0.824 A/W at 1.55 μm is comparable to literature values, although slightly lower than the 1.08 A/W reported for pn-based devices. Overall, the results indicate that the proposed nBn design achieves an excellent trade-off between ultra-low dark current and competitive optical performance.

4. Conclusion

In this work, an InGaAs-based nBn photodetector incorporating an $\text{Al}_{0.57}\text{Ga}_{0.43}\text{As}$ barrier layer was systematically designed and numerically analyzed using SILVACO ATLAS. Through detailed band alignment engineering, the optimal barrier composition was identified to achieve a near-zero valence band offset ($\Delta E_v \approx 35$ meV) while maintaining a sufficiently large conduction band offset ($\Delta E_c \approx 1.26$ eV). This configuration effectively enables barrier-free transport of minority carriers while strongly suppressing majority-carrier leakage, addressing one of the fundamental limitations of conventional InGaAs photodiodes. The optimized device demonstrates ultra-low dark current on the order of ~ 1 pA at low reverse bias (-0.3 V), representing orders-of-magnitude improvement compared to previously reported nBn-InGaAs structures. This significant reduction is primarily attributed to the suppression of Shockley–Read–Hall generation and diffusion-related leakage mechanisms, as confirmed by current component analysis. At higher biases, the controlled contribution of tunneling current further verifies the effectiveness of the barrier design in limiting field-assisted leakage. In addition to its low-noise performance, the proposed structure maintains competitive optical characteristics. By integrating a $\text{Si}_3\text{N}_4/\text{SiO}_2$ anti-reflection coating, surface reflectance was reduced to nearly zero at 1.55 μm , resulting in a responsivity enhancement of approximately 36% (from 0.60 A/W to 0.82 A/W). The device also exhibits low capacitance (~ 9.2 pF at -5 V) and a high breakdown voltage (~ 82.5 V), indicating its suitability for high-speed, low-noise, and wide dynamic range applications. Overall, this study demonstrates that precise control of barrier composition, doping, and layer thickness enables simultaneous optimization of electrical and optical performance in nBn architectures [26]. The presented design achieves an excellent trade-off between ultra-low dark current and high responsivity under low-bias operation, making it a strong candidate for next-generation short-wave infrared photodetectors operating under low-photon-flux and high-sensitivity conditions.

Funding

This research received no external funding.

Conflicts of interest

The authors declare that there are no conflicts of interest related to this article.

Data availability statement

The data is available from the corresponding author under reasonable request.

Author contribution statement

Ç. Tok was responsible for the conceptualization, design, simulation, formal analysis, investigation, methodology, resources, data curation, visualization, and writing of the original draft, as well as the review and editing of the manuscript.

M. Satılmış focused on the visualization, resources, methodology, investigation, and critical review of the manuscript

H. Keleş handled the data curation, methodology, investigation, and re-editing of the manuscript.

F. Oğuz provided conceptualization, supervision, and contributed to the re-editing and design.

H. Sarı provided supervision and contributed to the conceptual framework.

E. Özbay oversaw the methodology, conceptual framework, and final manuscript approval.

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